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# Skyrmions in van der Waals centrosymmetric materials

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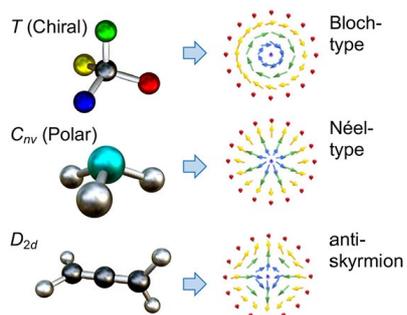
## Introduction

### Noncentrosymmetric vs Centrosymmetric

#### Noncentrosymmetric magnets

**Key magnetic interaction:**  
Dzyaloshinskii-Moriya interaction  
due to lack of space inversion symmetry

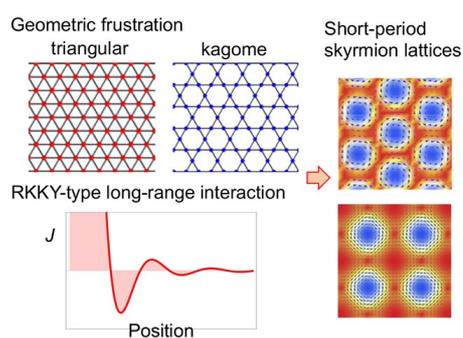
**Properties of skyrmion:**  
Fixed vorticity and helicity



#### Centrosymmetric magnets

**Key magnetic interactions:**  
Magnetic frustration and/or Higher-order interactions  
due to magnetic mediation by itinerant electrons

**Properties of skyrmion:**  
Ultra-small size (~ a few nm)



**Searching new mechanism to stabilize skyrmions in centrosymmetric lattice is urgent task**

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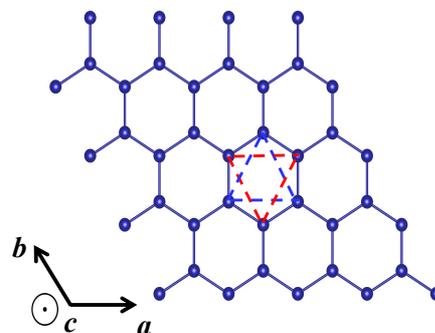
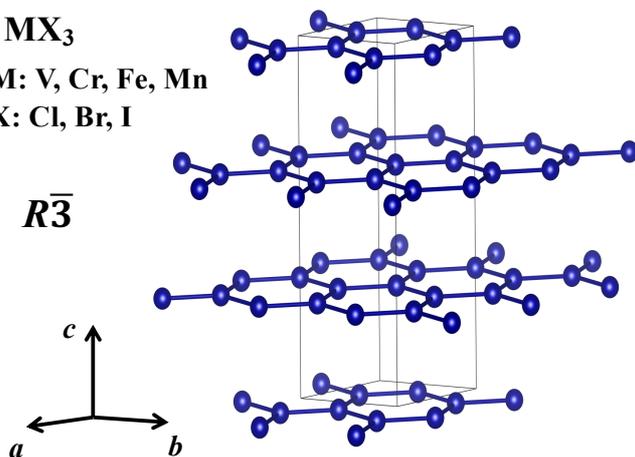
## Introduction Target materials

$\text{MX}_3$

M: V, Cr, Fe, Mn

X: Cl, Br, I

$R\bar{3}$



Dzyaloshinskii-Moriya interactions of 2nd NN can be finite

**Purpose: To investigate the effect of Dzyaloshinskii-Moriya interactions on magnetic properties of van der Waals centrosymmetric  $\text{MX}_3$**

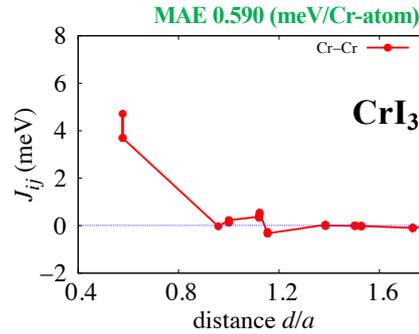
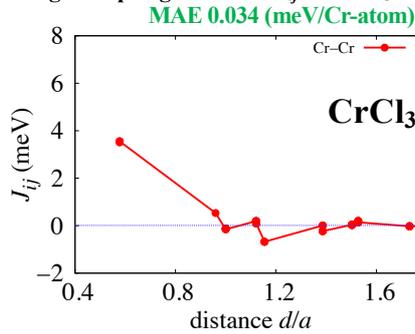
## Methods

### Classical Heisenberg model with Dzyaloshinski-Moriya interaction

$$H_{\text{Heis}} = - \underbrace{\sum_{\langle ij \rangle} J_{ij}^m \vec{S}_i \cdot \vec{S}_j}_{\text{Isotropic exchange}} - \underbrace{\sum_{\langle ij \rangle} \vec{D}_{ij} \cdot (\vec{S}_i \times \vec{S}_j)}_{\text{Antisymmetric exchange}} - \underbrace{\sum_i k_u (\vec{e}_u \cdot \vec{S}_i)^2}_{\text{Uniaxial anisotropy}} - \underbrace{g \mu_B \sum_i \vec{H}_{\text{ext}} \cdot \vec{S}_i}_{\text{Zeeman}}$$

## Results and Discussions

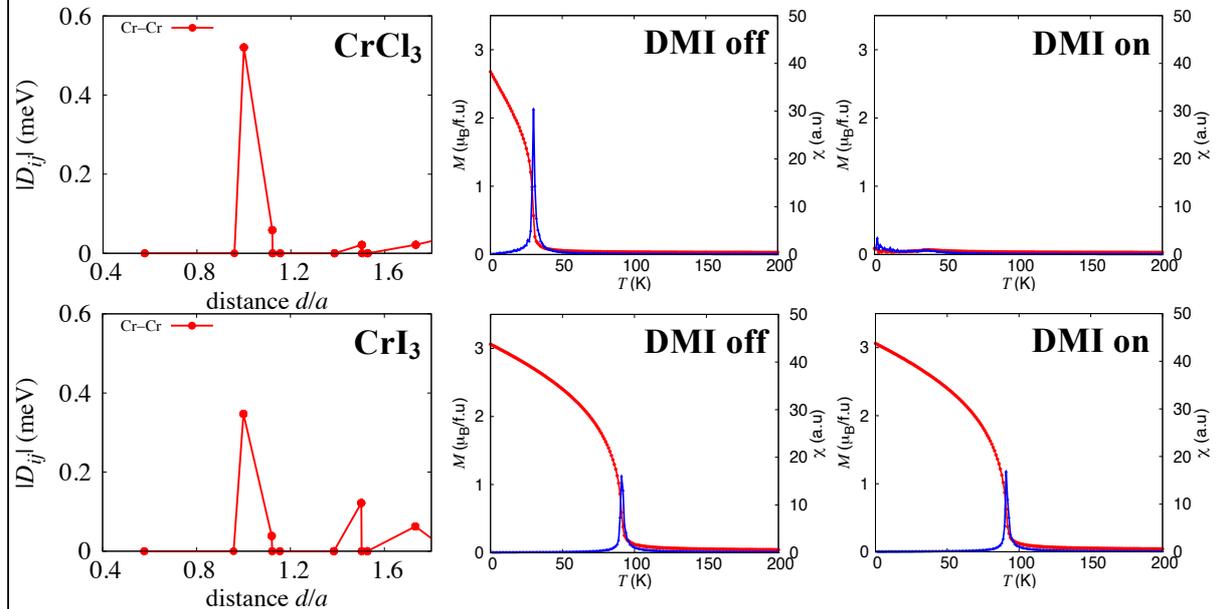
### Isotropic exchange coupling constant $J_{ij}$ of $\text{CrX}_3$



We use classical Heisenberg model with isotropic exchange, antisymmetric exchange or DMI, uniaxial anisotropy, and Zeeman term. All parameters of Monte Carlo simulations are obtained from first principles calculations.

First we consider the magnetic exchange coupling constants  $J_{ij}$  in  $\text{CrX}_3$ . When X change from Cl to I, the magnetocrystalline anisotropy increase from 0.034 to 0.590 meV/ Cr-atom. The positive value of MAE mean the easy axis is out of plane, the c direction. In additions, the interlayer exchange coupling constant in  $\text{CrCl}_3$  is positive. So it does not support the in plane AFM in experimental paper.

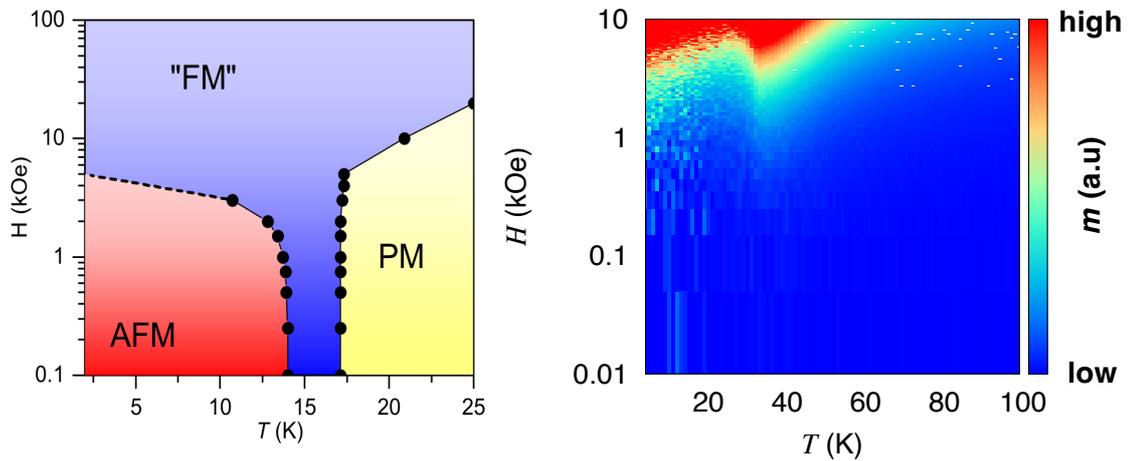
### $D_{ij}$ length and magnetization of $\text{CrX}_3$



Then we consider the DMI in first principles calculations. The DMI of second nearest neighbor of intralayer is finite and quite large even in  $\text{CrCl}_3$ . As point out in previous slide, the results of MAE and  $J_{ij}$  in  $\text{CrCl}_3$  does not support the inplane AFM in experimental works. And if we calculated the magnetization with switch off DMI, all  $\text{CrX}_3$  shows clear FM-PM transition with the ground state being FM.

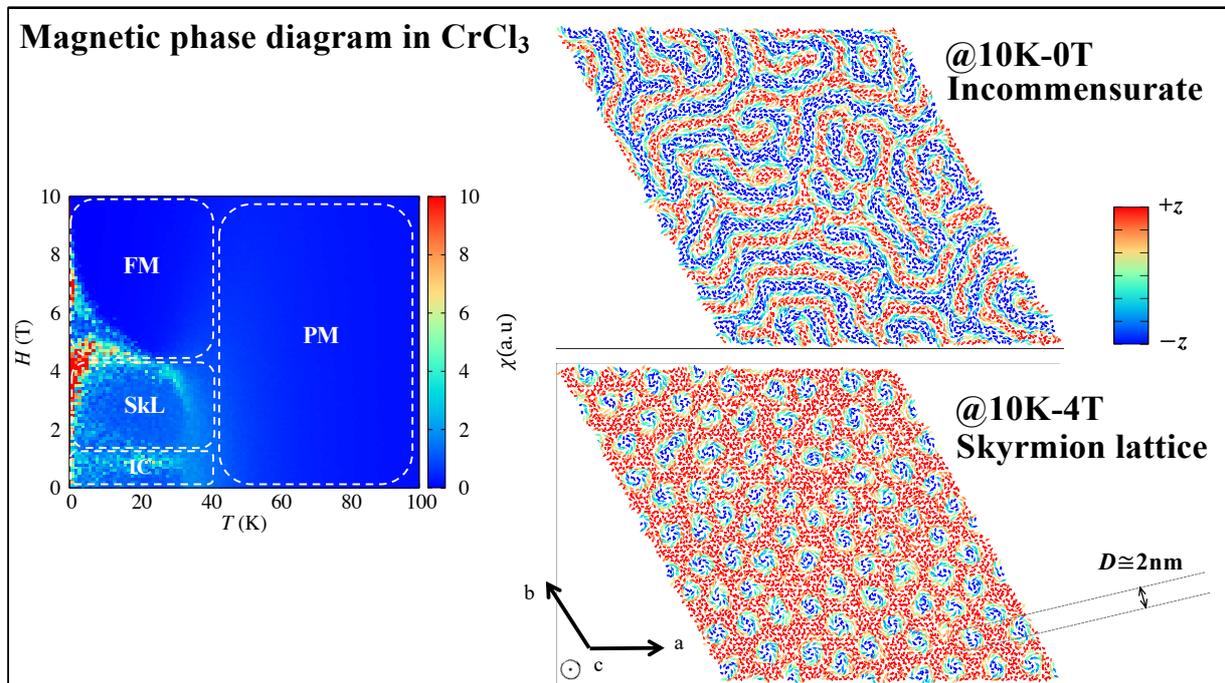
But if we include the DMI in simulations, the FM in  $\text{CrCl}_3$  is disappeared while the FM-PM transition is maintained in  $\text{CrBr}_3$  and  $\text{CrI}_3$  due to high MAE. So with the DMI in simulation, we can get zero magnetization and magnetic susceptibility in  $\text{CrCl}_3$ , similar behavior with experimental of inplane AFM order.

## Comparison of magnetic phase diagram of $\text{CrCl}_3$

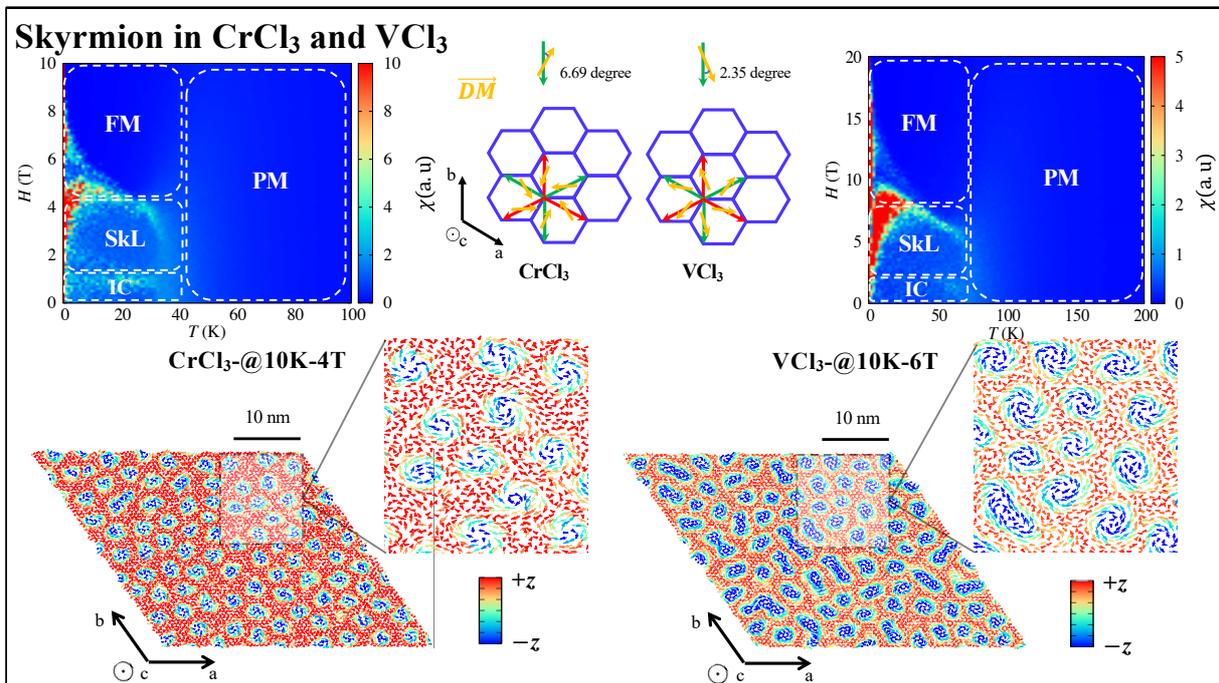


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Then we calculate the magnetic phase diagram of  $\text{CrCl}_3$ . the blue color means magnetization is small, it is AFM or PM in experimental phase diagram. The red color means the magnetization is high, it is FM in experimental. Here at this temperature, the red region is expand in low magnetic field region. We got a good agreement with experimetal study.

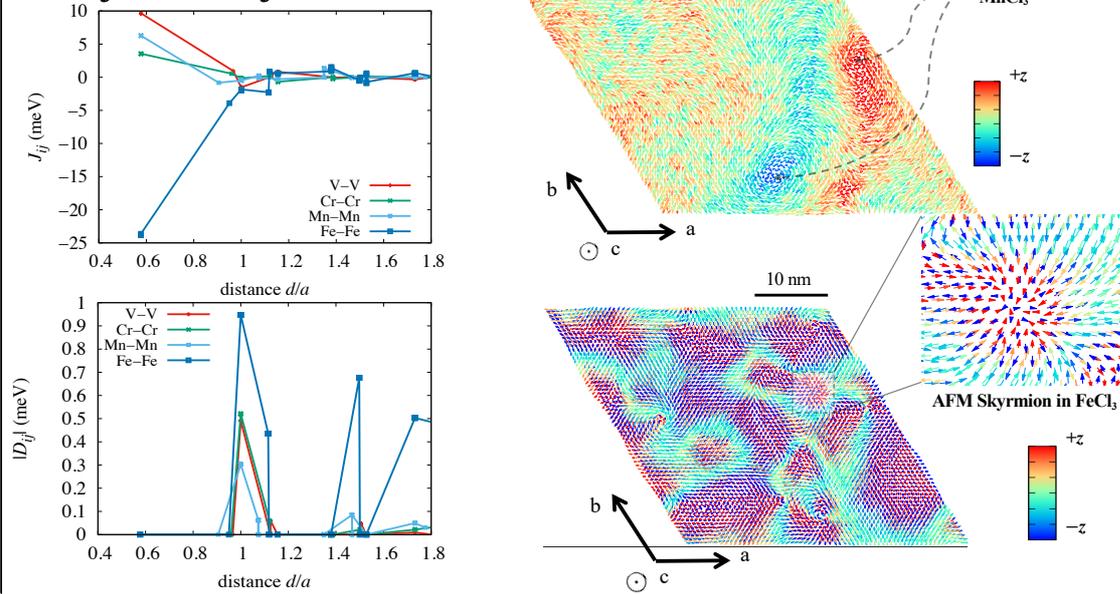


Then we check the snapshot of Monte Carlo simulations for making the magnetic phase diagram in magnetic susceptibility. At 10 K and 0 T, the magnetic phase is incommensurate with broken of long range order. If we apply magnetic field like 4 T, we can get the skyrmion lattice in  $\text{CrCl}_3$ .



We also found that  $\text{VCl}_3$  also have similar magnetic phase diagram as  $\text{CrCl}_3$ . However, the DM vector of  $\text{VCl}_3$  is nearly opposite with  $\text{CrCl}_3$ . It leads to two type of skyrmion with different helicity. If we can control the electron or hole dope in multilayer, we might control the helicity of the materials. Note that the skyrmion diameter is quite small, about 2 nm due to the centrosymmetric materials.

## Meron and AFM skyrmion in $\text{MnCl}_3$ and $\text{FeCl}_3$



We also found the Meron in  $\text{MnCl}_3$  and AFM skyrmion in  $\text{FeCl}_3$ . In  $\text{MnCl}_3$ , the magnetic anisotropy is easy plane with ab plane, it leads to the formation of Meron at zero external magnetic field.

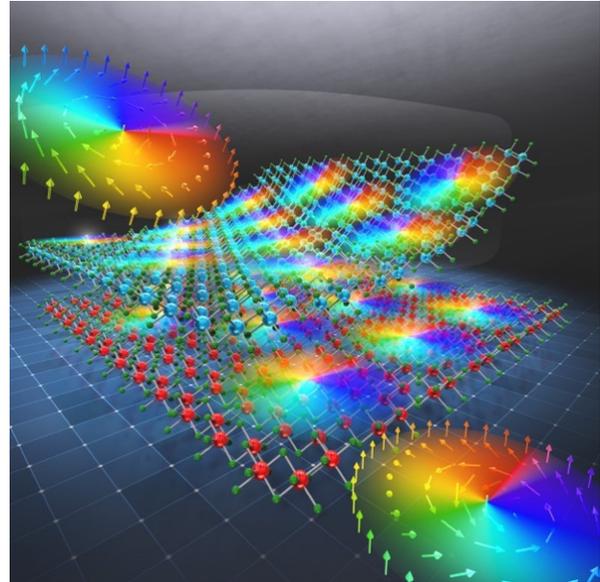
In  $\text{FeCl}_3$ , the magnetic exchange coupling constants is negative, so it prefer AFM order. The AFM skyrmion has advantage that the hall velocity is approximately equal to zero, so it can survivor for long time.

## Summary

**New mechanism to stabilize skyrmions in van der Waals centrosymmetric material**

**Not only skyrmions but also meron and antiferromagnetic skyrmions are observed**

**VCl<sub>3</sub> and CrCl<sub>3</sub> shows possibility to control helicity by heterostructure and electric field**



**H. B. Tran, Y. Matsushita, arXiv:2209.02333(2022)**

**Thank you so much for kind attentions!**

That all for our talk today, thank you so much for kind attentions.